

Molecular dynamics study of microstructural evolution and damage-depth induced by isotropic irradiation on Zr-1%Nb alloy

Abstract

Zirconium alloys with niobium have an extensive application in the nuclear industry, especially in fuel cladding. In this study, we consider the lattice structure of Zr-1%Nb alloy and study the damage depth (DD) due to irradiation on the structure of this alloy which results from the collision cascade (CC) phenomenon. It has been shown that the DD in the structure is directly related to Primary Knockedon Atom (PKA) energy. Because the structure of Zr-1%Nb is not homogeneous, DD is highly affected by the incident direction of irradiation. Both Zr and Nb atoms were considered as PKA's and the results show that the average of DD is larger for Nb than Zr. Next, the CC phenomenon has been considered for this alloy and the microstructure evolution has been studied at low temperature and low PKA energy. The results show the formation of some self-interstitials (SI's) during the CC phenomenon and no SI clusters have been observed.

Keywords: Zirconium-Niobium alloy, Collision-Cascade Phenomenon, Damage Depth, Molecular Statics

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INTRODUCTION

Because of low neutron absorption cross-section and high resistivity against corrosion, zirconium-niobium alloys are used as structural materials and fuel cladding in nuclear power reactors [1-4]. The study of microstructure evolution of the ZrNb alloy is, therefore, important when subjected to energetic irradiations due to fission in reactors.

According to available data, there are three phases for crystalline zirconium: α -phase ($< 863\text{ }^{\circ}\text{C}$), β -phase ($> 863\text{ }^{\circ}\text{C}$), and ω -phase (Pressure 2-7 GPa) [5-7]. Zr-Nb alloys are limited to low Nb percentages ($< 5\%$).

These alloys have a proper balance between corrosion resistance and low hydrogenation capacity [8]. In particular, Zr-1%Nb, which is studied in this work, has an HCP structure at low temperatures and niobium atoms are dissolved on Zr host structure. The structures of pure zirconium, niobium, and Zr-1%Nb alloy are shown in Figure 1. Under the irradiation, the atoms may displace irreversibly from their equilibrium lattice sites and give rise to the formation of defects. The resulted defects, then, may join together to make defect clusters. The macroscopic properties of nuclear materials are highly sensitive to the microstructural evolution of atoms in the materials and other parameters such as temperature and pressure. The mechanical properties of Zr-1%Nb alloy in presence of defects have already been studied by the present authors [9, 10]. Also, other studies have been done by others [7-8,11-20] and some aspects of this alloy and other similar materials have been considered until now. But there are still a lot of open questions that researchers are working on.

This work is divided into three main steps: First, lattice structure properties have calculated to verify the validity of the used interatomic potential in the calculations. Then, the damage depth (DD) value have calculated for Zr and Nb as PKA atoms , and the results of DD have compared. Finally, the CC phenomenon for low primary knocked-on atom (PKA) energies has investigated.

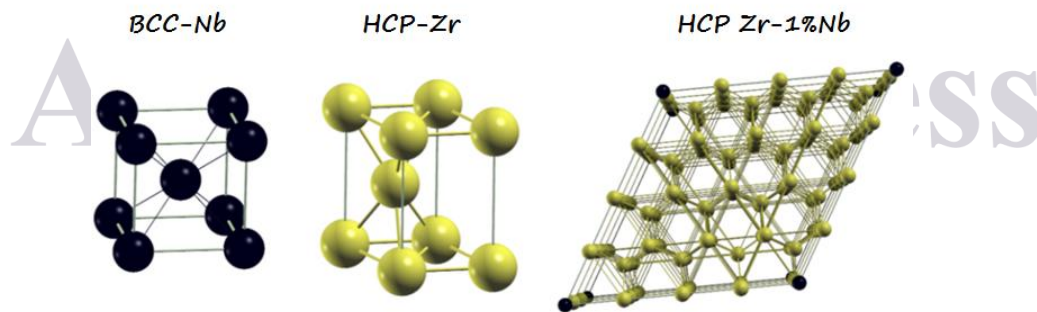


Fig1 Zirconium, Niobium and Zr-1%Nb alloy structures at low temperatures.

THEORIES AND COMPUTATIONAL DETAILS

All calculations have been done using atomistic simulation by the LAMMPS code package [21]. All post processing have been done by VMD [22] visualizer and gnuplot. Angular-dependent potential (ADP) is used as an interatomic potential for simulations. This potential is a new interatomic potential that is used to simulate Zr-Nb alloy and can predict its mechanical properties correctly [23]. This kind of interatomic potential is commonly used for metals and metal alloys and it is a generalization of EAM¹ potential that includes dipole and quadrupole distortions of the local environment of an atom. The simulation box for studying Zr-1%Nb is a 8×8×7 supercell that contains 43008 atoms (with 96 atoms per unit cell) with dimensions of 103.4×89.6×108.1 (°A)³. There is no any temperature effect in this study, so all calculations have been done at near zero (1 K) temperature. The structures have been fully optimized using the NPT ensemble and then the NVE ensemble has been used to simulate the CC phenomenon and DD processes. The 8 and 15 picoseconds were needed for equilibration of the structure in CC and DD phenomena when the structure has been subjected to radiation.

Table 1: Lattice constants in X & Z direction for pure-Zr (a,c) and Zr-1%Nb alloy in °A and VFE and SIFE for pure-Zr in eV. Reference data is composed of experimental and other available DFT calculations.

	Zr				Zr-1%Nb	
	<i>a</i>	<i>c</i>	<i>VFE</i>	<i>SIFE</i>	<i>a</i>	<i>c</i>
This work	3.188	5.206	1.9	2.76	12.856	15.360
Reference data	3.232[13]	5.149[13]	>1[24]	2.75[25]	-	-

RESULTS AND DISCUSSION

1. Lattice parameters

At first, the equilibrium lattice constants of pure-Zr in X and Z-directions (a,c) have been calculated and the results are in good agreement with experimental data. Because the structure of Zr-1%Nb is as pure-Zr, the calculations have been repeated to obtain the lattice constants of this alloy. Each unitcell of Zr-1%Nb has 96 atoms with an HCP structure. The results have been listed in Table 1. Vacancy and self-interstitial formation

¹ Embedded Atom Method

energies (VFE, SIFE) are two important parameters when one is considering the defects in the structure. In other words, the value of VFE and SIFE can show the behavior of the structure, when it is involved in the irradiation-induced defects. These quantities can be calculated as [7]:

$$E_V^f(p) = E(N - 1; P) - \frac{N - 1}{N} E(N; P) \quad (1)$$

$$E_{SI}^f(p) = E(N + 1; P) - \frac{N + 1}{N} E(N; P) \quad (2)$$

where $E_V^f(p)$, $E_{SI}^f(p)$ are VFE and SIFE in constant pressure respectively and N is the number of atoms in the structure. In addition, $E(N \pm 1; P)$, $E(N; P)$ are the total energy with one excess/less atom and perfect system. Formation energies for pure Zr structure have also been listed in Table 1 and compared with experimental results. A pair of vacancy and self-interstitial is called Frenkel-pair (FP) and the formation energy of FP's (FPFE) is simply the sum of VFE and SIFE and for pure-Zr FPFE is 4.66 eV . According to these results, the used interatomic potential is reliable to simulate the CC phenomenon.

2. Damage depth

The damage depth (DD) is the maximum distance of the PKA position to the furthest SI, that is, the maximum length which is affected by the irradiation. On the other hand, this factor is a damage criterion of the structure for each incident energetic particle. The DD of Zr-1%Nb alloy directly shows the strength of this alloy when it is subjected to irradiation and it is highly dependent on the irradiation with respect to crystalline orientations. This spatial angle dependency is depicted in Figures 2 and 3 for different PKA energies. X and Y and Z axis in Fig 2 are θ , ϕ and DD respectively. The orientation space of irradiation has meshed for different polar and azimuthal angles. The polar and azimuthal angles are varied in the ranges of (0-90°) and (0-120°) and the mesh step is 15° for each angle. The energy of PKA atoms in Fig2a to 2e are 10, 50, 100, 200 and 300 eV respectively.

When this alloy is subjected to irradiation, PKA atoms can be Zr or Nb atoms. Therefore, the results of damage of the structure can be different for different PKA atom types. In general (Fig2a to 2e), DD is bigger for Nb-PKA atoms in this structure. Therefore, the damage range for Nb as PKA atoms is more than Zr.

For $E(\text{PKA}) = 10$ eV, DD values are zero for nearly all directions. From experimental data, [26] the threshold displacement energy for Zirconium structure to produce a stable FP is about 22.5 eV. According to our calculation for PKA energy between 20-30 eV stable FP can be seen in most directions.

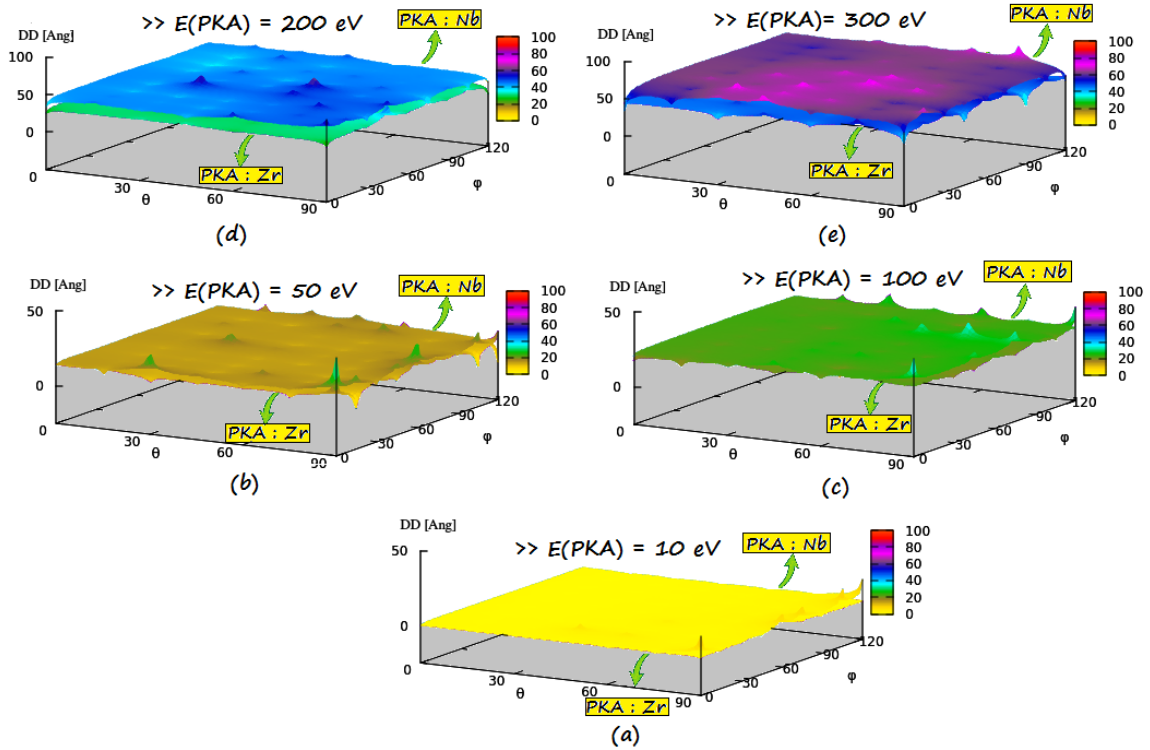


Fig 2 The dependency of DD to the polar and azimuthal angles of irradiation for different $E(\text{PKA})$: a) $E(\text{PKA}) = 10$ eV, b) $E(\text{PKA}) = 50$ eV, c) $E(\text{PKA}) = 100$ eV, d) $E(\text{PKA}) = 200$ eV, e) $E(\text{PKA}) = 300$ eV.

The left side of Figure 3 is related to Zr-PKA atoms and the right side to Nb. From this figure, it is clear that DD for $E(\text{PKA}) = 50$ eV for Zr and Nb-PKA are nearly the same and are changing $(0-40)^\circ\text{A}$. DD for $E(\text{PKA}) = 100, 200$ and 300 eV for Nb-PKA is meaningfully more than Zr-PKA and for 100 eV are in the range of $(11-37), (11-45)^\circ\text{A}$. This factor for 200 eV PKA energy is about $(16-62), (24-84)$. Finally, for 300 eV PKA energy, DD is changed as $(19-73)$ and $(34-90)^\circ\text{A}$ for Zr- and Nb-PKA atoms. It is shown that the maximum DD is happened for a head-on collision with nearest neighbor atoms such as X-direction ($\theta = 90$ & $\phi = 0$), as can be seen in Figure 3.

Furthermore, there are 8 SI positions in Zr-1%Nb which show in Figure 4 and the results show that for $E(\text{PKA}) = 300 \text{ eV}$, the maximum of DD can be nearly seen in the direction of the octahedral SI position ($\theta = 55.4^\circ$ & $\phi = 30^\circ$). In other words, when the direction of irradiation is in the SI direction of the structure, the furthest atoms can be placed in the SI positions after losing their energy during the CC process.

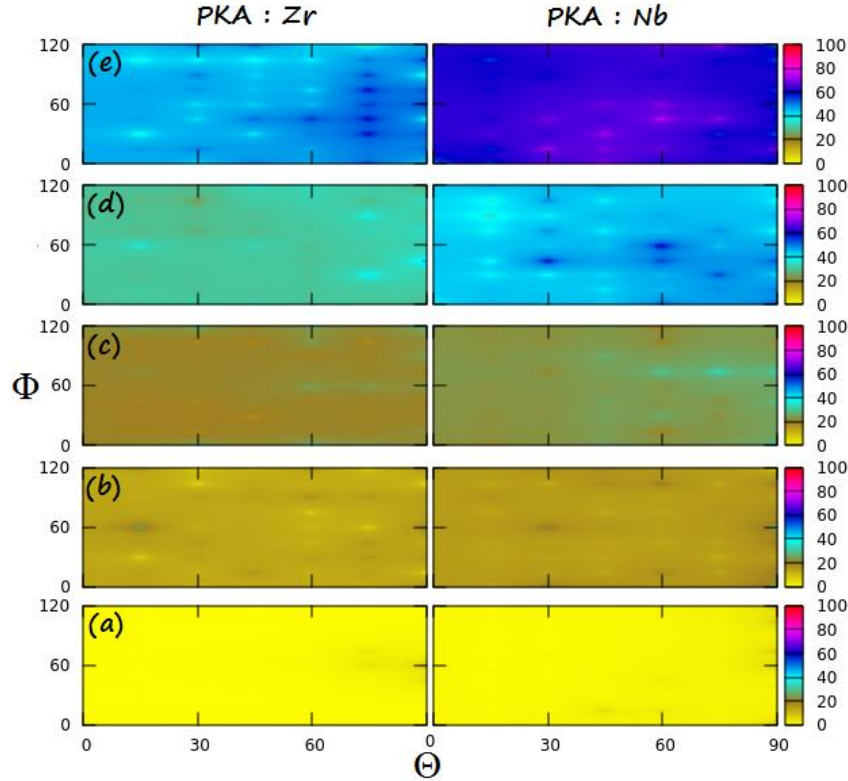


Fig 3 The 2D map dependency of DD to the polar and azimuthal angles of irradiation for different $E(\text{PKA})$: a) $E(\text{PKA}) = 10 \text{ eV}$, b) $E(\text{PKA}) = 50 \text{ eV}$, c) $E(\text{PKA}) = 100 \text{ eV}$, d) $E(\text{PKA}) = 200 \text{ eV}$, e) $E(\text{PKA}) = 300 \text{ eV}$. The first (V) and second vacancy positions (N1 to N5) in the HCP zirconium structure; N1 to N5 refer to 1st to 5th nearest neighbor structure positions.

Table 2: The average DD for different $E(\text{PKA})$ in $^\circ \text{A}$ and eV respectively.

$E(\text{PKA})$	10	50	100	200	300
AVGDD; Zr-PKA	0.8	13.7	21.4	31.5	47.2
AVGDD; Nb-PKA	1.7	16.0	24.1	45.2	63.2

It has also been shown that an increase in the $E(\text{PKA})$ can lead to an increase in DD. Furthermore, the treatment of the mean value of all DD in these orientations with respect

to the PKA energy have a deviation from linear relation and is shown in table 2 and Figure 5. In other words, the DD variation with respect to E(PKA) is similar to NRT model. When PKA energy for this alloy is near 300 eV , it means that the order of irradiation energy is about a few hundred eV . Hence, when the Zr-1%Nb structure is subjected to irradiation with relatively high energy the structure is affected nearly by $40 \sim 60 \text{ }^\circ\text{A}$.

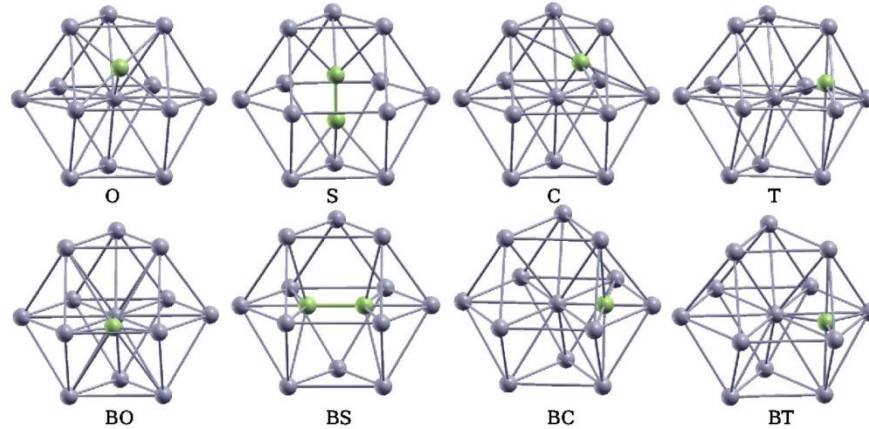


Fig 4 Eight SI positions in Zr-1%Nb alloy; O: octahedral, S: split-dumbble, C: crowdion, T: tetrahedral and B in 4 bottom stand for basal [7].

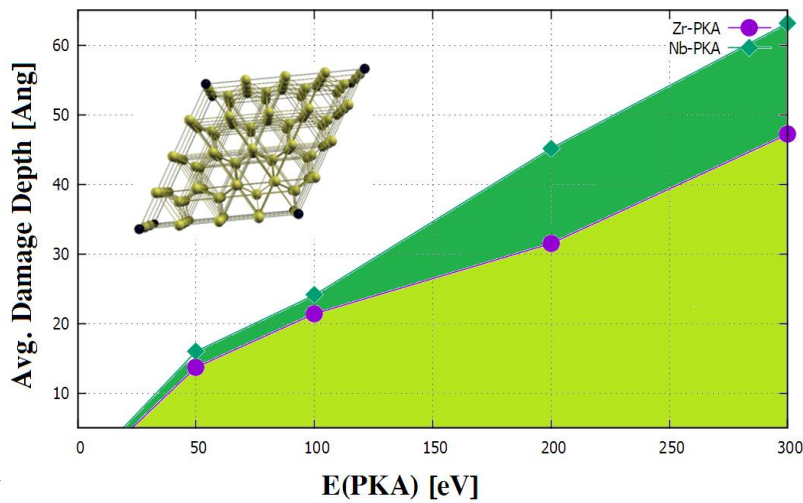


Fig 5 DD variation w.r.t. E(PKA) for Zr-1%Nb alloy.

3. Collision Cascade (CC) phenomenon

CC phenomenon has been considered for Zr-1%Nb alloy in this section. This phenomenon is illustrated in Figure 6. When PKA atoms have sufficient kinetic energy to transfer to the atoms nearby the PKA atoms, they are displaced from their equilibrium crystal lattice sites, leaving behind a vacant site or a vacancy and each displaced atom and its vacancy make an FP. The FP's, themselves, may have acquired sufficient energy to displace other atoms and this process can be repeated to make the CC phenomenon. At first, the number of FPs increase, then some of the displaced atoms recombine with vacant sites and the number of FPs decrease over time. So all of the available FPs are not stable and just a few FPs will remain in the structure. For Zr-1%Nb alloy, the most disturbed structure takes place around 0.5 picoseconds and the equilibration is achieved after 15 picoseconds. As it can be seen, no clustering of defects has been detected individually for low PKA energies and temperature in a perfect crystal. It is clear that for high PKA energies it is possible to find clusters of Self-Interstitials (SIs) or voids (V).

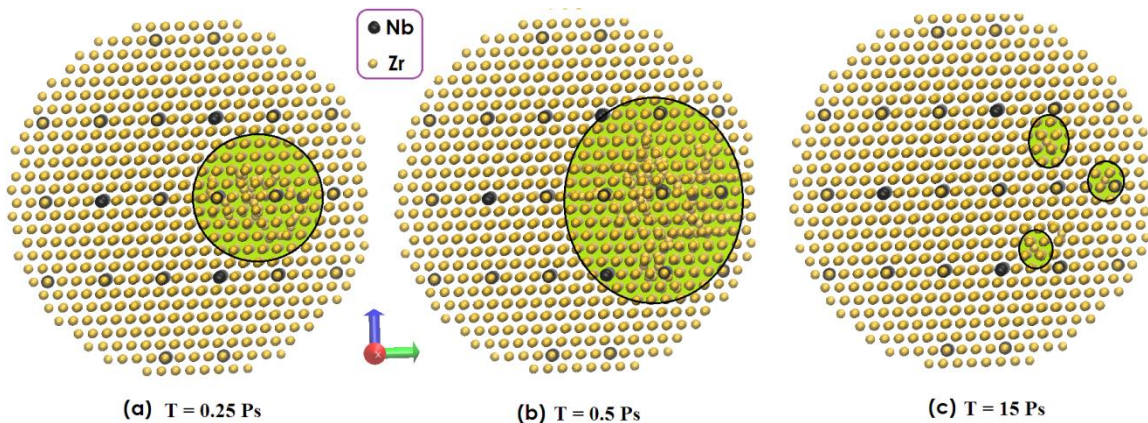


Fig 6 CC phenomenon for Zr-1%Nb alloy for different time steps.

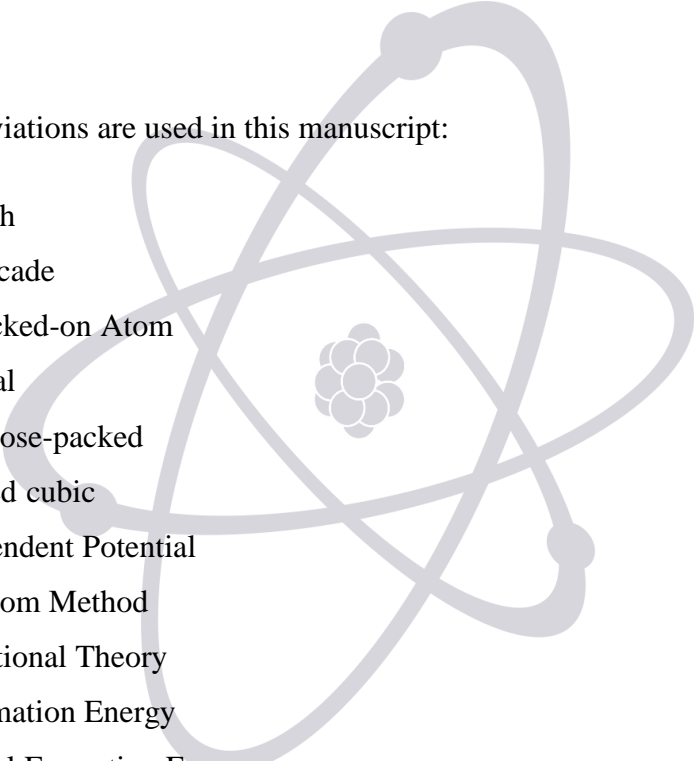
CONCLUSION

In this study, the DD of the PKA atoms with different energies for Zr-1%Nb alloy has calculated and the results have shown that the increase in PKA energy can lead to an increase in DD. In addition, the Nb and Zr atoms may be taken as PKA atoms and the change in the type of PKA atoms changes the damage depth of the alloy. It has been shown that Nb-PKA atoms can produce more damage to the structure. Furthermore, DD was meaningfully dependent on irradiation orientation and this dependency has been

explored in homogeneous angle meshes and it has been shown that in some cases, the direction of SI in HCP structure may lead to higher degrees of damage. The CC phenomenon was also considered for this alloy and it has been shown that a few Frenkel pairs were produced in low PKA energies.

Abbreviations

The following abbreviations are used in this manuscript:



DD	Damage Depth
CC	Collision Cascade
PKA	Primary Knocked-on Atom
SI	Self-Interstitial
HCP	Hexagonal Close-packed
BCC	Body-Centered cubic
ADP	Angular Dependent Potential
EAM	Embedded Atom Method
DFT	Density Functional Theory
VFE	Vacancy Formation Energy
SIFE	Self-Interstitial Formation Energy
FP	Frenkel Pair
FPFE	Frenkel Pair Formation Energy
NRT	Norgett, Robinson, and Torres

References

- Article in Press**
- [1] BDC Bell, ST Murphy, PA Burr, RJ Comstock, JM Partezana, RW Grimes and MR Wenman. *Corrosion Science* **105** PP. 36–43, 2016.
 - [2] G Fernandez and et. al. *Zeitschrift für Metallkunde* **82** (6) PP. 478–487, 1991.
 - [3] B Lustman and F Kerze, The metallurgy of zirconium. (McGraw-Hill Book Company 1955 4).
 - [4] B Cox, *Journal of Nuclear Materials* **336** (2) P 331–368, 2005.
 - [5] H Xia, S J Duclos, A L Ruoff and Y K Vohra, *Phys. Rev. Lett.* **64** (2), 1990.
 - [6] V O Kharchenko and D O Kharchenko, arXiv preprint arXiv 1206.7035, 2012.

- [7] Q Peng, W Ji, J Lian, X J Chen, H Huang, F Gao and S De, *Scientific reports*, Nature Publishing Group **41** (1), P 1-7, 2014.
- [8] D Smirnova and S Starikov, Defect and Diffusion Forum, *Trans Tech Publ.* **375**, P 167–174, 2017.
- [9] M R Basaadat and M Payami, *International Journal of Modern Physics C* **31** (02), P 2050028, 2020.
- [10] M R Basaadat and M Payami, *IRANIAN JOURNAL OF PHYSICS RESEARCH* **20** (1), P 57–64, 2020.
- [11] X K Xin, W S Lai and B X Liu, *Journal of Nuclear Materials* **393** (1), P 197–202, 2009.
- [12] X Wang, L B Liu, M F Wang, X Shi, G X Huang and L G Zhang, *Calphad* **48** (1), P 89–94, 2015.
- [13] ES Fisher and CJ Renken, *Physical review* **135** (2A), P A482, 1964.
- [14] J Ashkenazi, M Dacorogna, M Peter, Y Talmor, E Walker and S Steinemann, *Physical Review B* **18** (8), P 4120, 1978.
- [15] DI Bolef, *Journal of Applied Physics* **32** (1), P 100–105, 1961.
- [16] P F Weck, E Kim, V Tikare and J A Mitchell, *Dalton Transactions* **44** (43), P 18769–18779, 2015.
- [17] DJ Hayes and FR Brotzen, *Journal of Applied Physics* **45** (4), P 1721–1725, 1974.
- [18] YP Varshni, *Physical Review B* **2** (10), P 3952, 1970.
- [19] L Fast, JM Wills, B Johansson and O Eriksson, *Physical Review B* **51** (24), P 17431, 1995.
- [20] N Al-Zoubi, S Chonecker, X Li, W Li, B Johansson and L Vitos, *Computational materials science* **159**, P 273–280, 2019.
- [21] S J Plimpton LAMMPS features and capabilities. Sandia National Lab.(SNL-NM), Albuquerque, NM (United States) 2014.
- [22] W. Humphrey, A. Dalke, and K. Schulten, "VMD - Visual Molecular Dynamics", *J. Molec. Graphics*, vol. **14**, pp. 33-38, 1996.
- [23] DE Smirnova and SV Starikov, *Computational Materials Science* **129**, P 259–272, 2017.
- [24] ML Swanson, GR Piercy, GV Kidson and AF Quenneville, Atomic Energy of Canada Ltd., Chalk River (Ontario). Chalk River Nuclear Labs 1969.
- [25] G D Samolyuk, S I Golubov, YN Osetsky and R E Stoller. *Philosophical magazine letters* **93** (2), P 93–100, 2013.
- [26] WG Wolfer, *Comprehensive nuclear materials 1*, P 1–45, 2012.